where $\theta = \varepsilon_x + \varepsilon_y + \varepsilon_z$. In uniaxial strain, $\varepsilon_x = (V_0 - V)/V_0$, $\varepsilon_y = \varepsilon_z = 0$. If the equality holds, the material is in the plastic state.

ii) In the plastic state, every increment in strain is the sum of an elastic and a plastic increment:

(43*a*)
$$\mathrm{d}\varepsilon_x = \mathrm{d}\varepsilon_x^e + \mathrm{d}\varepsilon_x^p \,,$$

(43b)
$$\mathrm{d}\varepsilon_{y} = \mathrm{d}\varepsilon_{y}^{e} + \mathrm{d}\varepsilon_{y}^{p},$$

(43c)
$$\mathrm{d}\varepsilon_z = \mathrm{d}\varepsilon_z^e + \mathrm{d}\varepsilon_z^p \,.$$

iii) There is no plastic dilatation:

(44)
$$\mathrm{d}\varepsilon_x^p + \mathrm{d}\varepsilon_y^p + \mathrm{d}\varepsilon_z^p = 0 \; .$$

iv) The stress is supported solely by the elastic strain:

(45a)
$$\mathrm{d}p_x = \lambda \,\mathrm{d}\theta + 2\mu \,\mathrm{d}\varepsilon_x^e,$$

(45b)
$$\mathrm{d}p_{v} = \lambda \,\mathrm{d}\theta + 2\mu \,\mathrm{d}\varepsilon_{v}^{e} \,,$$

(45c)
$$\mathrm{d}p_z = \lambda \,\mathrm{d}\theta + 2\mu \,\mathrm{d}\varepsilon_z^{*}\,,$$

where λ and μ are, in general, functions of the density.

As p_x is increased from zero, the response is initially elastic and $\varepsilon_y = \varepsilon_z = 0$. Then

(46)
$$p_x - p_v = (1 - 2v) p_x / (1 - v)$$
,

where $v = \lambda/2(\lambda + \mu)$ is Poisson's ratio. The yield stress is reached at a value of p_x called the «Hugoniot elastic limit», denoted by p_{HEL} . From eqs. (41) and (46):

(47)
$$p_{\text{HEL}} = (1-\nu) Y/(1-2\nu)$$
.

For further increases in p_x , the material is in the plastic state. Then

(48)
$$p_x \equiv \overline{p} + \frac{2}{3}(p_x - p_y) = \overline{p} + 2Y/3,$$

where $\overline{p} = (p_x + p_y + p_z)/3$, a function of density and internal energy alone. Referring to Fig. 14 b), eq. (48) applies to the segment *AB* of the p_x curve. The slope of the (p_x, V) curve in the elastic region is, from eqs. (42):

(49)
$$dp_x/dV = -(\lambda + 2\mu)/V_0 = -(K + 4\mu/3)/V_0,$$

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where K is bulk modulus. In the plastic region, AB, the slope is, for constant Y, from eq. (48)

(50)
$$\mathrm{d}p_x/\mathrm{d}V = \mathrm{d}\overline{p}/\mathrm{d}V = -K/V.$$

In accord with eq. (50), it is convenient to define the incremental dilatation as dV/V. Bulk modulus normally increases with \overline{p} , so AB is normally concave upward. The yield stress, Y, is in general a function of plastic work and density. In such case eq. (50) is augmented by a dY/dV term. In any case the offset of p_x from the hydrostat, \overline{p} , is always 2Y/3.

At point B in Fig. 14 b) we suppose that a change is made from monotonically increasing to monotonically decreasing p_x . Equation (41) must again be examined to determine whether the mass element is in the elastic or plastic state. During the initial compression process, p_x increased more rapidly than p_y until yield occurred. During unloading, p_x decreases more rapidly than p_y until yielding again occurs. Thus the portion BC of the unloading curve is elastic until $p_y - p_x = Y$ at C. From C to D, unloading is plastic and the unloading curve lies below the hydrostat by $\frac{2}{3}Y$.

Referring to the discussion following eq. (17), we see that point A of Fig. 14 b) may be a point of instability for single shock compressions. To see that this is indeed the case, suppose that a shock wave has been generated with amplitude p_{HEL} , traveling with speed

$$D_{\scriptscriptstyle E} = [V_{\scriptscriptstyle 0}(\lambda + 2\mu)]^{\frac{1}{2}}.$$

The velocity of this shock front relative to the material behind it is

(51)
$$D_{e} - u_{E} = (V_{A}/V_{0}) D_{E} = V_{A} \sqrt{(\lambda + 2\mu)/V_{0}}$$

If an additional compression of small amplitude is produced to follow the already established shock, it will travel with velocity c_{A} relative to the material ahead of it, where, according to eq. (50),

$$c_{A} = \sqrt{KV_{A}} = V_{A}\sqrt{(\lambda + 2\mu/3)/V_{A}}.$$

Comparing this with eq. (51) we find that

(52)
$$(D_{\mathcal{B}} - u_{\mathcal{B}})^2 / c_{\mathcal{A}}^2 = (3V_{\mathcal{A}}/V_0)(1-\nu)/(1+\nu) \simeq 3(1-\nu)/(1+\nu) = \frac{3}{2} \text{ for } \nu = \frac{1}{3},$$

since $V_A/V_0 \simeq 1$ at the Hugoniot elastic limit. According to eq. (52), the second wave does not overtake the shock, so there is a region of the (p_x, V) curve above the point A which cannot be reached by a single shock from

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